**Quick Summary** 

$$mu'' + \gamma u' + ku = F(t)$$

Case 1: F(t) = 0 and 
$$\gamma = 0$$
  
 $u(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t)$   
 $\omega_0 = \sqrt{k/m} = natural freq.$ 

**Case 2**: F(t) = 0 and 
$$\gamma > 0$$
  
**2a**:  $\gamma > 2\sqrt{mk}$ , **overdamped**  
**2b**:  $\gamma = 2\sqrt{mk}$ , **critically damped**  
**2c**:  $\gamma < 2\sqrt{mk}$ , **damped vibrations**  
 $u(t) = e^{\lambda t} (c_1 \cos(\mu t) + c_2 \sin(\mu t))$   
 $\mu = \sqrt{\frac{k}{m} - \frac{\gamma^2}{4m^2}} =$  **quasi**-frequency

Case 4:  $F(t) = F_0 \cos(\omega t)$  and  $\gamma > 0$ Sol'n:  $u(t) = u_c(t) + u_p(t)$  $u_c(t)$  =homogeneous sol'n = transient sol'n  $u_p(t)$  =particular sol'n = steady state sol'n (also called forced response)

We will discuss Case 4 today.
Here is an example:
Entry Task:
m = 1 kg, γ = 2 N/(m/s), k = 5 N/m.
External Forcing: F<sub>0</sub> = 10 N, ω = 1 rad/s.
1.Find the homogenous solution.
2.Find a particular solution.

**Case 3**:  $F(t) = F_0 \cos(\omega t)$  and  $\gamma = 0$   $u(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + u_p(t)$  **3a**:  $\omega \neq \omega_0$ :  $u_p(t) = \frac{F_0}{m(\omega_0^2 - \omega^2)} \cos(\omega t)$  **3b**:  $\omega = \omega_0$ :  $u_p(t) = \frac{F_0}{2m\omega_0} t \sin(\omega t)$ Resonance!

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The general solution to

u'' + 2u' + 5u = 10\cos(t)

is

u(t) =

e^{-t}(c_1\cos(2t) + c_2\sin(2t)) + 2\cos(t) + \sin(t)
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*Example:* Same problem as before with smaller damping:

m = 1 kg,  $\gamma$  = 0.1 N/(m/s), k = 5 N/m. External Forcing: F<sub>0</sub> = 10 N,  $\omega$  = 1 rad/s.  $u'' + 0.1u' + 5u = 10\cos(t)$ 

Solution:

 $u(t) = e^{-0.05t} (c_1 \cos(\mu t) + c_2 \sin(\mu t))$  $+ 2.4984 \cos(t) + 0.0624 \sin(t)$ Here is the graph with c<sub>1</sub> = 4 and c<sub>2</sub> = -6.



## Some First Observations:

- 1. When there is damping, there is a *transient* part of the solution that always dies out.
- 2. If damping is smaller, it takes longer to die out.
- 3. The amplitude of the steady state solution is dependent on m,  $\gamma$ , k, and F<sub>0</sub> in some way.

Consider the example

 $u'' + \gamma u' + 5u = 10\cos(\omega t)$ 

Homogenous Solution: The characteristic equation is  $r^2 + \gamma r + 5 = 0$   $r = -\frac{\gamma}{2} \pm \frac{1}{2}\sqrt{\gamma^2 - 20}$ If  $\gamma < \sqrt{20}$ , then  $\mu = \frac{1}{2}\sqrt{20 - \gamma^2} = \sqrt{5 - \frac{\gamma^2}{4}}$ Note:  $w_0 = \sqrt{5}$ 

Particular Solution: Using:  $u_p(t) = A\cos(\omega t) + B\sin(\omega t)$ we get

$$\begin{array}{l} (5-\omega^2)A+\gamma\omega B=10\\ -\gamma\omega A+(5-\omega^2)B=0 \end{array}$$

$$A = \frac{10(5 - \omega^2)}{\omega^4 + (\gamma^2 - 10)\omega^2 + 25}$$
$$B = \frac{10\gamma\omega}{\omega^4 + (\gamma^2 - 10)\omega^2 + 25}$$

as you can see it starts to get messy.

Let's look at the case when  $\omega = \sqrt{5}$ then A = 0 and B =  $\frac{10\gamma\sqrt{5}}{25+(\gamma^2-10)5+25} = \frac{2\sqrt{5}}{\gamma}$  $u_p(t) = \frac{2\sqrt{5}}{\nu} \sin(\sqrt{5}t)$ When  $\gamma = 0.1$ , you get

$$u'' + \gamma u' + 5u = 10\cos(\sqrt{5} t)$$

Steady state solution:

$$u_p(t) = \frac{2\sqrt{5}}{\gamma} \sin(\sqrt{5} t)$$

γ	R
10	0.447
1	4.47
0.1	44.72
0.01	447.21
0.001	4472.14

## Some Second Observations:

- 1. If the forcing frequency is close to the natural frequency, then tend to get large amplitude solutions.
- 2. In this case, the amplitude gets larger and larger the closer the damping is to zero.

## **General Discussion**

$$mu'' + \gamma u' + ku = F_0 \cos(\omega t)$$

Note: 
$$\omega_0 = \sqrt{\frac{k}{m}}$$
 ,  $\mu = \sqrt{\frac{k}{m} - \frac{\gamma^2}{4m^2}}$ 

Particular Solution:

 $u_p(t) = A\cos(\omega t) + B\sin(\omega t)$ 

Leads to

$$-\gamma \omega A + (k - m\omega^2)B = 0$$
$$(k - m\omega^2)A + \gamma \omega B = F_0$$

The formulas for A and B are large to write out.

The amplitude of the steady state solution simplifies to:

$$R = \sqrt{A^{2} + B^{2}} = \frac{F_{0}}{\sqrt{(k - m\omega^{2})^{2} + \gamma^{2}\omega^{2}}}$$

Thinking of this as a function of  $\omega$  the maximum steady state amplitude occurs when:

$$\omega_{max} = \sqrt{\frac{k}{m} - \frac{\gamma^2}{2m^2}}$$

In particular, for small values of  $\gamma$ if  $\omega \approx \omega_0$ , then  $R \approx \frac{F_0}{\gamma \omega}$  is large. (resonance)